

欧拉 - 泊松方程组的自相似解*

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摘要: 三维可压等熵欧拉泊松方程组描述可压等熵理想气态星体的运动规律, 它由质量守恒方程、动量守恒方程及自引力位势满足的泊松方程构成。研究欧拉泊松方程组的自相似解是天体物理及数学领域研究的热点问题之一, 具有重要的现实意义和广阔的应用前景。关于三维可压等熵欧拉泊松方程组自相似解的研究较少。用分离变量法研究了一类三维可压等熵欧拉泊松方程组的一组自相似解。

关键词: 三维可压等熵; 欧拉泊松方程组; 自相似解

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Self-similar solutions for the compressible Euler-Poisson equations in three dimension

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Abstract: The isentropic compressible Euler-Poisson equations, addressed to describe the motion of ideal gaseous stars, consist of the Euler-Poisson equations for the conservation of mass and momentum, and Poisson equation induced by the potential function of the self-gravitational force. Using the separation method, self-similar solution are presented for the isentropic compressible Euler-Poisson equations in three dimension.

Key words: isentropic compressible; Euler-Poisson equations; self-similar solutions; three dimension

在天体物理学中, 三维可压等熵欧拉 - 泊松方程组为如下形式:

$$\begin{cases} p_t + \nabla \cdot (pv) = 0 \\ p[v_t + (v \cdot \nabla)v] + k \nabla p^\gamma = -p \nabla \varphi \end{cases} \quad (1)$$

和

$$\Delta \varphi = 4\pi gp \quad (2)$$

其中 $x = (x_1, x_2, x_3) \in \mathbf{R}^3$; $p = p(t, x)$ 和 $v = v(t, x) = (v_1, v_2, v_3) \in \mathbf{R}^3$, $\varphi = \varphi(t, x)$ 分别代表流体的

密度, 运动速度和自引力位势, g 是引力常数, 为了计算简便, 设 $g = 1, k > 0; \gamma \gg 1$ 是常数。

1 概述

欧拉 - 泊松方程组不仅在天体物理学有着重要的应用, 而且在生物科学等许多领域亦有着广泛的应用。它主要研究受到电场外力或自引力作用的流体的运动。在流体动力学模型中, 欧拉 - 泊松方程

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组描述半导体器件或是等离子体中电子的运动；在生物学模型中，描述蛋白质通道中离子的传送，即离子在细胞外部和薄膜的细胞质边之间的传动^[1]。三维欧拉 - 泊松方程组是描述，具有自引力势能的气态星体内部结构发展的流体动力学的一个模型，对 Euler-Poisson 方程组解的研究，是天体物理学家和许多数学学者研究的热点问题。文献 [2] 给出该方程组静止状态下，关于解的存在性和不存在性的一些结果；文献 [3] 给出了方程组在某种情况下会有爆破解的结果；文献 [4] 给出了方程组平衡解的存在性及多重性；文献 [5] 给出了方程组 $1 < \gamma < \frac{6}{5}$ 时，可压欧拉 - 泊松方程组平衡解的存在性。本文研究三维可压等熵 Euler-Poisson 方程组，它由质量守恒方程，动量守恒方程及自引力位势满足的泊松方程构成。文献 [6] 研究了该方程组一类爆破解和时间周期解；文献 [7] 研究了该方程组的平衡解的存在唯一性；文献 [8] 研究了该方程组平衡解的稳定性，而关于三维可压等熵欧拉泊松方程组自相似解的研究较少。对于欧拉泊松方程组的自相似解的研究，也是天体物理及数学领域研究的热点问题之一，具有重要的现实意义和广阔的应用前景。

关于非线性偏微分方程组自相似解的求法，郭志荣等^[9]研究了运用相似变换的复合变换法将高维非线性偏微分方程组化为常微分方程，从而求出高维非线性偏微分方程组的自相似解；赵向青等^[10]中提到了 Ribaud 等通过选择一类特殊的初始值函数而获得了经典 Schrödinger 方程的自相似解。而求解偏微分方程的精确解有各种各样的方法：傅立叶变换法、格林函数法、逆散射方法、Fokas 变换法、分离变量法等^[11]。由文献 [12] 得到启发，我们研究了一类三维可压等熵欧拉泊松方程组，用分离变量法得到了该方程组的自相似解。

2 主要结论及其证明

定理 1 三维可压等熵欧拉 - 泊松方程组，存在下面一组自相似解：

$$\begin{cases} p = \frac{f(s)}{a^2(t)b(t)}, \\ v_1 = \frac{\dot{a}(t)}{a(t)}x_1 - \frac{\tau}{a^2(t)}x_2, \\ v_2 = \frac{\tau}{a^2(t)}x_1 + \frac{\dot{a}(t)}{a(t)}x_2, \\ v_3 = \frac{\dot{b}(t)}{b(t)}x_3 \end{cases} \quad (3)$$

这里自相似变量 $s = \frac{x_1^2 + x_2^2}{a^2(t)} + \frac{x_3^2}{b^2(t)}$ 而

$$f(s) = \left(\alpha^{\gamma-1} - \frac{\beta}{2k\gamma}s - \frac{G(t)}{2k\gamma} \right) \int_0^s \frac{1}{x} \int_{\mathbb{R}^3} p(y,t) \frac{x-y}{|x-y|} dy ds \Big)^{\frac{1}{\gamma-1}} \quad (\gamma > 1) \quad (4)$$

与式 (4) 对应的有

$$\begin{cases} \ddot{a}(t) - \frac{\tau^2}{a^3(t)} = \frac{\beta}{a^{2\gamma-1}(t)b^{\gamma-1}(t)}, \\ a(0) = a_0 > 0, \dot{a}(0) = a_1, \\ G(t) = a^{2\gamma}(t)b^{\gamma-1}(t); \\ \ddot{b}(t) = \frac{\beta}{a^{2\gamma-2}(t)b^\gamma(t)}, \\ b(0) = b_0 > 0, \dot{b}(0) = b_1, \\ G(t) = a^{2\gamma}(t)b^{\gamma+1}(t) \end{cases} \quad (5)$$

这里 $\tau \neq 0, \beta, \alpha \geq 0, a_0, a_1, b_0, b_1$ 是任意常数

证明 文献 [12] 已证自相似解 (3) 满足质量守恒方程 $p_t + \nabla \cdot (pv) = 0$ 。下面证明自相似解 (3) 满足方程组 (1) 的第二个方程。

方程组 (1) 的第二个方程的第一个分量方程是

$$p \{ v_{1t} + (v_1 v_{1x_1} + v_2 v_{1x_2} + v_3 v_{1x_3}) \} + k \frac{\partial p^\gamma}{\partial x_1} = -p \frac{\partial \varphi}{\partial x_1} \quad (6)$$

将自相似解 (3) 代入式 (6)，整理得

$$\begin{aligned} & \frac{p}{a^{2\gamma-1}(t)b^{\gamma-1}(t)a(t)} \frac{x_1}{a(t)} \cdot \\ & \left\{ \left(\ddot{a}(t) - \frac{\tau^2}{a^3(t)} \right) a^{2\gamma-1}(t)b^{\gamma-1}(t) + \right. \\ & \left. 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma}(t)b^{\gamma-1}(t) \frac{1}{x_1} \frac{\partial \varphi}{\partial x_1} \right\} = \\ & \frac{p}{a^{2\gamma-1}(t)b^{\gamma-1}(t)a(t)} \frac{x_1}{a(t)} \cdot \\ & \left\{ \beta + 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma}(t)b^{\gamma-1}(t) \frac{1}{x_1} \frac{\partial \varphi}{\partial x_1} \right\} = 0 \end{aligned}$$

其中

$$\begin{cases} \ddot{a}(t) - \frac{\tau^2}{a^3(t)} = \frac{\beta}{a^{2\gamma-1}(t)b^{\gamma-1}(t)}, \\ a(0) = a_0 > 0, \dot{a}(0) = a_1, \\ \beta + 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma}(t)b^{\gamma-1}(t) \frac{1}{x_1} \frac{\partial \varphi}{\partial x_1} = 0 \end{cases} \quad (7)$$

且

$$f(0) = \alpha \geq 0, G(t) = a^{2\gamma}(t)b^{\gamma-1}(t) \quad (8)$$

由方程 (2) 得

$$\varphi(x) = - \int_{\mathbb{R}^3} \frac{p(y,t)}{|x-y|} dy \quad (9)$$

则可求得方程 (8) 的精确解是

$$f(s) = \left(\alpha^{\gamma-1} - \frac{\beta}{2k\gamma}s - \frac{G(t)}{2k\gamma} \right).$$

$$\int_0^s \frac{1}{x_1} \int_{\mathbb{R}^3} p(y, t) \frac{x_1 - y_1}{|x - y|} dy ds \Big)^{\frac{1}{\gamma-1}} (\gamma > 1) \quad (10)$$

方程组 (1) 的第二个方程的第二个分量方程是

$$p \{ v_{2t} + (v_1 v_{2x_1} + v_2 v_{2x_2} + v_3 v_{2x_3}) \} + k \frac{\partial p^\gamma}{\partial x_2} = -p \frac{\partial \varphi}{\partial x_2} \quad (11)$$

将自相似解 (3) 代入式 (11), 整理得

$$\frac{p}{a^{2\gamma-1}(t)b^{\gamma-1}(t)} \frac{x_2}{a(t)} \cdot$$

$$\left\{ \left(\ddot{a}(t) - \frac{\tau^2}{a^3(t)} \right) a^{2\gamma-1}(t)b^{\gamma-1}(t) + \right.$$

$$\left. 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma}(t)b^{\gamma-1}(t) \frac{1}{x_2} \frac{\partial \varphi}{\partial x_2} \right\} =$$

$$\frac{p}{a^{2\gamma-1}(t)b^{\gamma-1}(t)} \frac{x_2}{a(t)} \cdot$$

$$\left\{ \beta + 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma}(t)b^{\gamma-1}(t) \frac{1}{x_2} \frac{\partial \varphi}{\partial x_2} \right\} = 0$$

其中

$$\ddot{a}(t) - \frac{\tau^2}{a^3(t)} = \frac{\beta}{a^{2\gamma-1}(t)b^{\gamma-1}(t)},$$

$$a(0) = a_0, \dot{a}(0) = a_1,$$

$$\beta + 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma}(t)b^{\gamma-1}(t) \frac{1}{x_2} \frac{\partial \varphi}{\partial x_2} = 0$$

且

$$f(0) = \alpha \geq 0, G(t) = a^{2\gamma}(t)b^{\gamma-1}(t) \quad (12)$$

则可求得方程 (12) 的精确解是

$$f(s) = \left(\alpha^{\gamma-1} - \frac{\beta}{2k\gamma}s - \frac{G(t)}{2k\gamma} \right).$$

$$\int_0^s \frac{1}{x_2} \int_{\mathbb{R}^3} p(y, t) \frac{x_2 - y_2}{|x - y|} dy ds \Big)^{\frac{1}{\gamma-1}} (\gamma > 1) \quad (13)$$

方程组 (1) 的第二个方程的第三个分量方程是:

$$p \{ v_{3t} + (v_1 v_{3x_1} + v_2 v_{3x_2} + v_3 v_{3x_3}) \} + k \frac{\partial p^\gamma}{\partial x_3} = -p \frac{\partial \varphi}{\partial x_3} \quad (14)$$

将自相似解 (3) 代入式 (14), 整理得

$$\frac{p}{a^{2\gamma-1}(t)b^{\gamma+1}(t)} \frac{x_3}{a(t)} \left\{ \ddot{b}(t)a^{2\gamma-2}(t)b^\gamma(t) + \right.$$

$$\left. 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma-1}(t)b^{\gamma+1}(t) \frac{1}{x_3} \frac{\partial \varphi}{\partial x_3} \right\} = 0$$

其中

$$\begin{cases} \ddot{b}(t) = \frac{\beta}{a^{2\gamma-2}(t)b^\gamma(t)}, \\ b(0) = b_0 > 0, \dot{b}(0) = b_1 \end{cases}$$

$$\beta + 2k\gamma f^{\gamma-2}(s)f(s) + a^{2\gamma-1}(t)b^{\gamma+1}(t) \frac{1}{x_3} \frac{\partial \varphi}{\partial x_3} = 0$$

且

$$f(0) = \alpha \geq 0 \quad G(t) = a^{2\gamma}(t)b^{\gamma+1}(t) \quad (15)$$

则可求得方程 (15) 的精确解是

$$f(s) = \left(\alpha^{\gamma-1} - \frac{\beta}{2k\gamma}s - \frac{G(t)}{2k\gamma} \right).$$

$$\int_0^s \frac{1}{x_3} \int_{\mathbb{R}^3} p(y, t) \frac{x_3 - y_3}{|x - y|} dy ds \Big)^{\frac{1}{\gamma-1}} (\gamma > 1) \quad (16)$$

这就证明了三维可压等熵欧拉泊松方程组有一组自相似解, 定理 1 证毕。

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